

Low (and high) temperature plasma Laser Diagnostics

Alexandros Gerakis

Princeton Plasma Physics Laboratory Introduction to Fusion Energy and Plasma Physics Course June 15^{th} , 2023











General Questions about Luxembourg

- It is <u>not</u> in Germany
- Yes, it is its own country.
- No, it is <u>not</u> a kingdom but a Grand Duchy (the last one in the world).
- One of the founding members of the EU.
- It has three(!) official languages:
 - □English is <u>not</u> one of them
 - ✓ Luxembourgish (yes, it is a thing)
 - ✓ German
 - \checkmark French
- Population is roughly 615.000
- Almost same size as Rhode Island
- 260 times smaller than Texas.



THE NEW RESEARCH AND TECHNOLOGY ORGANIZATION



Created in 2015 Active in 4 fields

- ENVIRONMENT
- ► INFORMATICS
- ► MATERIALS
- ► SPACE





ABOUT US FUNDING INNOVATION & INDUSTRY PARTNERSHIPS INTERNATIONAL COOPERATION SCIENCE IN SOCIETY ISC "Photonics &

Home | Our News | Dr Alexandros Gerakis awarded FNR ATTRACT Fellowship

Dr Alexandros Gerakis awarded FNR ATTRACT Fellowship

18 Feb 2022

ON MA PH · / · ORATOR

WHO I AM

G f У 🖶 🕇

The FNR has awarded a FNR ATTRACT Fellowship to Dr Alexandros Gerakis for an innovative laser technology project at the Diving deep into the theory of plasma Institute of Science and Technology (LIST). Dr Gerakis' AT doted with 2 MEUR over five years.

Single-shot coherent

physics

~3

NDING

oelectronic Devices" HERIOT

WATT

St Andrews

UNIVERSITY University of

August 8, 2019 | By Jan McHarg

Aerospace Engineering National Science Foundation (NSF) Research



Drs. Alexandros Gerakis and Kentaro Hara | Image: Igor Kraguljac

With a \$300,000 grant from the National Science Foundation, a team of researchers at Texas A&M University is diving deep into the physics of plasma to question the fundamental theory of local thermodynamic equilibrium (LTE) in arc discharges.

ACKNOWLEDGEMENTS

Optical Probing & Manipulation Group







Dr. Junhwi Bak Post-Doctoral Research Associate

Robert Randolph Graduate Research Associate

Dr. Mikhail N. Shneider



Dr. Kentaro Hara Stanford University



TEXAS A&M





Office of Science



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ATTRACT

ATTRACTING OUTSTANDING YOUNG RESEARCHERS FROM ABROAD





DISCLAIMER NO.1

This is not an exhaustive list

Research Capabilities



Advanced Diagnostics

- · Atomic and molecular species, positive and negative ions
 - Laser-Induced Fluorescence (LIF) and Two-Photon Absorption LIF (TALIF)
 - Nanosecond LIF/TALIF Femtosecond TALIF
 - CW LIF for measurements of Ar and Xe ion and atom VDFs
- ant Enhanced Multi-Photon Ionization (Ra
- Characterization of plasma, chemical composition and dynamic behavior
- Optical Emission Spectro
- Fourier Transform Infrared Spect
- drupole Mass Spectro meter and Energy Analyzer
- Electron velocity distribution function, gas flow velocity and temperature
- rent Anti-Stokes Raman Scattering (Hybrid CARS)
- osecond Laser Electronic Excitation Tagging (FLEET)
- Electric field and space potential
- Electric Field-Induced Second Harmonic Generation (E-FISH)
- Nanoparticle Diagnostics
- nduced Breakdown Sp
- avleigh-Brillouin Scatterin
- Surface Diagnostics
 - ond Harmonic Generation (SHG)
 - n-induced Secondary Electron Emission and Surface Charging
 - Surface Potential Measurements with the Kelvin Probe Diag

https://pcrf.princeton.edu/capabilities/#diagnostics

Capabilities

- Multiple femtosecond, picosecond, nanosecond and CW lasers spanning < 190 nm and exceeding 10 mm for interrogating dynamic plasma and reactive environments.
 - LCIF for electron densities and LIF-Dip for electric fields.
 - 2D-CARS and SFG for gas and surface phase interrogation
 - Multi-pulse (10 kHz to 100 kHz) laser cluster for high speed LIF and Rayleigh
- High speed imaging capabilities from <200 ps gated cameras, multi-frame framing cameras and high-speed CMOS cameras.
- VUV and FTIR spectrometers for spectroscopy.
- High resolution tandem and molecular beam mass spectroscopy for interrogating gas phase chemistries occurring in multi-atmosphere.

Broad range of plasma generating capabilities spanning vacuum to atmospheric pressures and beyond.

https://www.sandia.gov/prf/plasma-research-facility/capabilities/



B. Laser-Collision-Induced Fluorescence (LCIF)

Pure He 2.3% H₂O shroud

P2:10.11ac

Installered

Quartz frit, sample packed

ES 16-2 June





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DISCLAIMER NO.2

I will be going fast ©



WHY CARE ABOUT METROLOGY?

and Light

<u>Metrology:</u> Tools used to measure/quantify a process

Metrology is defined by the International Bureau of Weights and Measures (BIPM) as "the science of measurement, embracing both experimental and theoretical determinations at any level of uncertainty in any field of science and technology".



From https://www.energy.gov/ne/articles/role modeling-and-simulation-scientific-discovery

WHY CARE ABOUT LTP METROLOGY?



Hall Thruster



Nishida, Hiroyuki, Taku Nonomura, and Takashi Abe. "Three-dimensional simulations of discharge plasma evolution on a dielectric barrier discharge plasma actuator." *Journal of Applied Physics* 115, no. 13 (2014): 133301.



Arc Jets https://www.nasa.gov/centers/ames/multi media/images/2008/A-28916-3.html



Magnetron sputtering https://intlvac.com/News-Resources/ArticleID/5/Wha

esources/ArticleID/5/What-is-Sputtering



S. Yatom, A. Khrabry, J. Mitrani, A. Khodak, I. Kaganovich, V. Vekselman,
B. Stratton and Y. Raitses "Synthesis of nanoparticles in carbon arc: measurements and modeling", MRS. Comm. (2018)





M. Kundrapu, M. Keidar, Numerical simulation of carbon arc discharge for nanoparticle synthesis, Physics of Plasma, Vol. 19, No. 073510. 2012

 Important quantities

 for every flow:

 • Temperature

 • Density

 • Flow velocity

 • Flow velocity

 • WHAT EVEN IS

 TEMPERATURE????



WHY CARE ABOUT FLOW METROLOGY?

Ideal characteristics for **every** diagnostic:

✓ Nonperturbative





 \checkmark Single

✓ Works with all ✓ Works *in-situ*, in a multitude of environments (e.g. wind-tunnels/fire/plasma)











WHY CARE ABOUT FLOW METROLOGY? The state-of-the-art



Measurement of <u>flow velocity with the Pitot</u> <u>tube</u>, improved by M. H. Darcy (Recherches hydrauliques, entreprises par M. H. Darcy, continuées par M. H. Bazin. Extract from volume XIX of Mémoires présentés par divers savants à L'institut Impérial de France, 1st part (501 p., 28 pl.): Recherches expérimentales sur l'écoulement de l'eau dans les canaux découverts. Imprimerie impériale, Paris 1865, planche IV).









WHY CARE ABOUT FLOW METROLOGY? The state-of-the-art

Plasma & Plasma flow diagnostics: Langmuir Probes



 $Taken \ from \ https://www.youtube.com/watch?v{=}6uZfcE80DDQ$



Hot Wire



WHY CARE ABOUT FLOW METROLOGY? The solution....

https://www.eurekalert.org/multimedia/697040





"LASERS"

https://phys.org/news/2021-11-laser-cooling-quantum-gases.html



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Light **Amplification** (by) Stimulated Emission (of) Radiation





Fifty years ago, five scientists from Culham Laboratory made a trip to the Soviet Union that was to prove a pivotal moment in the quest for fusion energy.

Their journey to Moscow's Kurchatov Institute in 1969 confirmed the impressive results of the Soviets' T3 experiment – a relatively new type of fusion device known as a 'tokamak'. The Culham team measured plasma temperatures of around 10 million degrees C at T3; far in excess of any previous fusion machine and a big step towards the conditions needed for fusion power. The emergence of the tokamak was a huge boost for fusion research and opened the way to machines like JET, ITER and a viable route to future power plants.

These days, international collaboration in science is taken for granted. But in the late 1960s, the Cold War was at its peak and mutual suspicion between East and West made the Culham scientists' mission an almost unprecedented breach of the 'Iron Curtain'.

To mark the 50th anniversary of the publication of the T3 results in Nature, we spoke to Dr Mike Forrest, who was on the team that went to Moscow. Mike helped to develop the **'Thomson scattering' laser techniques** that made the measurements possible and that are still a key diagnostic for today's tokamaks. Mike recalls the pioneering early use of lasers, the top-level clearance needed for the trip and how UK and Soviet scientists bonded over a shared love of tea as they made a breakthrough that still resonates today.

50 years on: The mission to Moscow that changed fusion research



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https://ccfe.ukaea.uk/mission-to-moscow-50-years-on/

WITHOUT LASERS.... Optical Emission Spectroscopy







https://www.ugent.be/ea/appliedphysics/en/research/plasma/plasmaphysicsandengineering.htm/plasmadiagnostics.htm and the second second

Dictionary

Definitions from Oxford Languages 🞯 · Learn more 🥝



noun

the degree or intensity of heat present in a substance or object, especially as expressed according to a comparative scale and shown by a <u>thermometer</u> or perceived by touch. "at a temperature of 2°C"

- the degree of internal heat of a person's body.
 "I'll take her temperature"
- INFORMAL
 a body temperature above the normal.
 "he was running a temperature"

, Wikipedia

https://en.wikipedia.org > wiki > Temperature

Temperature - Wikipedia

Temperature is a physical quantity that expresses quantitatively the perceptions of hotness and coldness. Temperature is measured with a thermometer. Thermodynamic temperature $\bigcirc \chi freeting \cdot Atmospheric temperature \bigcirc \chi freeting \cdot Scale of temperature \odot \chi freeting \cdot Scale of tempe$

> LUXEMBOURG INSTITUTE OF SCIENCE AND TECHNOLOGY

Temperature is a measure of the average kinetic energy per particle in a substance.







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https://commons.wikimedia.org/wiki/File:Carbon_monoxide_rotational-vibrational_spectrum.png

When:

Translational =:: Vibrational =:: Rotational =:: Electronic

Thermodynamic Equilibrium



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SO....WHAT IS TEMPERATURE?

When:

Translational # Vibrational # Rotational # Electronic

Thermodynamic Non-Equilibrium



Temperature is defined through a Maxwellian distribution function...

What if our particles do <u>**not**</u> follow a Maxwellian distribution function?



How is temperature defined then???



WHAT DO WE WANT TO MEASURE IN A PLASMA???



Xenon (@15 Torr)/ Gerakis lab





Footage from the inside of the NSTX-U during a run. Image Courtesy of PPPL.



WHAT DO WE WANT TO MEASURE IN A PLASMA???



Electrons: Density, Temperature, Flow velocity

Atoms (neutrals): Density, Temperature, Flow velocity

Atomic ions: Density, Temperature, Flow velocity

Molecules (neutrals): Density, Temperature, Flow velocity



Temperature

Vibrational Temperature

Molecular ions: Density, Temperature, Flow velocity



HOW DO WE MEASURE THINGS IN A PLASMA???



HOW DO WE MEASURE THINGS IN A PLASMA???

LINEAR REGIME



THOMSON SCATTERING





Electronic temperature



THOMSON SCATTERING • EIG

Electrons



https://www.ipp.cas.cz/vedecka_struktura_ufp/tokamak/COMPASS/diagnostics/spektroskopicke-diagnostiky/thomsonuv-rozptyl.html

Typical cross-sections: $10^{-29} - 10^{-30} \text{ m}^{-2}$









Image taken from https://ceej.travellerspoint.com/157/





Atomic ions

Molecules (neutrals)









Image taken by AG somewhere in the Aegean




https://www.ugent.be/ea/appliedphysics/en/research/plasma/plasmaphysicsandengineering.htm/plasmadiagnostics.htm



RESONANT & NON-RESONANT APPROACHES



(elastic)



(inelastic)



SPONTANEOUS RAMAN SCATTERING





DOI 10.1088/0957-0233/12/5/201





Atoms (neutrals)

Atomic ions

Molecules (neutrals) Molecular ions



https://researchrepository.wvu.edu/cgi/viewcontent.cgi?article=6578&context=etd





LASER INDUCED FLUORESCENCE

Tunable laser Light sheet optics Inaging optics

https://andor.oxinst.com/learning/view/article/planar-laser-induced-fluorescence-as-a-plasma-diagnostic



a) Front b) Tail Reactor nozzle -Vertical 10 position (mm) 12-14 16-18. 20 22

2D TALIF Image of O atom: ns vs fs comparison

Atoms (neutrals)

Atomic ions

[Schmidt et al., Plasma Sources Science and Technology 26 (5), 055004 (2017)]



Molecules (neutrals)

Comparison of 2D image of the O atom employing (a) ns-TALIF and (b) fs-TALIF for the same discharge in 2%-O2/He-mixture APPJ. Minimum detection limit of fs-TALIF is factor-of-three lower than ns-TALIF.

H-atom plasma image

[Schmidt et al., Journal of Physics D: Applied Physics 50 (1), 015204 (2017)]





Molecular ions

https://spectral energies.com/research/electric-field-plasmas/plasma-imaging/

CAPPS/Gerakis lab

 $\bullet 10.1088/1402 - 4896/acc906$

LINEAR PLASMA DIAGNOSTICS WORK GREAT! Until....



Other neutral and plasma flow laser diagnostics, resonant and non-resonant, e.g. vibrationally excited nitric oxide monitoring (VENOM), femtosecond laser electronic excitation tagging (FLEET) laser induced fluorescence (LIF), Coherent Anti-Stokes Raman Scattering (CARS) and others have also been demonstrated to measure flow LUXEMBOURG Velocity and/or temperature.



HOW DO WE MEASURE THINGS IN A PLASMA???

NON-LINEAR REGIME





As the light intensity increases, this *linear* model begins to breakdown. Instead, we now use a power series to model the polarization of the material :

$$\boldsymbol{P_{ind}} = \varepsilon_0(\chi^{(1)} \cdot \boldsymbol{E} + \chi^{(2)} \cdot \boldsymbol{E}^2 + \chi^{(3)} \cdot \boldsymbol{E}^3 + \dots)$$

[†]Griffiths, D. J., (2012) Introduction to Electrodynamics, 4th Edition; Pearson Boyd, R.W., (2008) Nonlinear Optics, 3rd Edition; Elsevier



As the light intensity increases, this *linear* model begins to breakdown. Instead, we now use a power series to model the polarization of the material:

$$\boldsymbol{P_{ind}} = \varepsilon_0(\chi^{(1)} \cdot \boldsymbol{E} + \chi^{(2)} \cdot \boldsymbol{E}^2 + \chi^{(3)} \cdot \boldsymbol{E}^3 + \dots)$$

Considering an electric field of the form:

$$\boldsymbol{E}(t) = E_{\omega} \cos(\omega t)$$

We obtain for the induced polarizability:

$$\begin{aligned} \boldsymbol{P}_{ind} &= \varepsilon_0(\boldsymbol{\chi}^{(1)} \cdot \boldsymbol{E}_{\omega} \cos(\omega t) + \boldsymbol{\chi}^{(2)} \cdot \boldsymbol{E}_{\omega}^2 \cos^2(\omega t) + \boldsymbol{\chi}^{(3)} \cdot \boldsymbol{E}_{\omega}^3 \cos^3(\omega t) + \dots) \\ & \text{Non-centrosymmetric} \\ & \text{materials } (\boldsymbol{\chi}^{(3)} = 0) \\ \boldsymbol{P}_{ind} &= \varepsilon_0(\boldsymbol{\chi}^{(1)} \cdot \boldsymbol{E}_{\omega} \cos(\omega t) + \frac{1}{2}\boldsymbol{\chi}^{(2)} \cdot \boldsymbol{E}_{\omega}^2 + \frac{1}{2}\boldsymbol{\chi}^{(2)} \cdot \boldsymbol{E}_{\omega}^2 \cos(2\omega t)) \\ & \text{Term} \\ & \text{oscillating at} \\ & \text{frequency } \omega \\ & \text{(linear)} \\ & \text{propagation)} \\ \end{aligned}$$

As the light intensity increases, this *linear* model begins to breakdown. Instead, we now use a power series to model the polarization of the material:

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As the light intensity increases, this *linear* model begins to breakdown. Instead, we now use a power series to model the polarization of the material:

$$\boldsymbol{P_{ind}} = \varepsilon_0(\chi^{(1)} \cdot \boldsymbol{E} + \chi^{(2)} \cdot \boldsymbol{E}^2 + \chi^{(3)} \cdot \boldsymbol{E}^3 + \dots)$$

Implies that three fields interact

to produce a fourth field. Hence,

 $\chi^{(3)}$ interactions are four

photon processes!

Considering an electric field of the form:

$$\boldsymbol{E}(t) = E_{\omega} \cos(\omega t)$$

We obtain for the induced polarizability:

• We know that, for the electric displacement, we have for a linear system:

$$D = \varepsilon_0 \cdot E + P_{ind}$$

= $\varepsilon_0 \cdot E + \varepsilon_0 \cdot \chi^{(1)} \cdot E$
= $\varepsilon_0 \cdot (1 + \chi^{(1)}) \cdot E$
= $\varepsilon_0 \cdot \varepsilon_r \cdot E$

where ε_r is the relative permittivity of the medium.

• We define the refractive index of the medium as:

$$n = \sqrt{\varepsilon_r} = \sqrt{1 + \chi^{(1)}}$$



• Considering a centro-symmetric medium (i.e. $\chi^{(2)}=0)$ the polarization of the material will be:

$$P(\omega) = \varepsilon_0 \cdot \chi^{(1)} E(\omega) + \varepsilon_0 \cdot \chi^{(3)} E^3(\omega)$$

= $\varepsilon_0 \cdot E(\omega)(\chi^{(1)} + \chi^{(3)} E^2(\omega))$

• Hence for the electric displacement we will have:

$$\boldsymbol{D}(\omega) = \boldsymbol{\varepsilon}_0 \cdot \boldsymbol{E}(\omega) + \boldsymbol{\varepsilon}_0 \cdot \boldsymbol{E}(\omega)(1 + \chi^{(1)} + \chi^{(3)}\boldsymbol{E}^2(\omega))$$

= $\boldsymbol{\varepsilon}_0 \cdot \boldsymbol{E}(\omega)(1 + \chi^{(1)} + \chi^{(3)}\boldsymbol{E}^2(\omega))$

• Noting that $E^2(\omega)$ is simply the intensity $I(\omega)$ of the beam, we can express the relative permittivity of the medium as:

$$\varepsilon_r = (1 + \chi^{(1)} + \chi^{(3)}I(\omega))$$

• This yields for the refractive index of the medium:

$$n = \sqrt{\varepsilon_r} = \sqrt{(1 + \chi^{(1)} + \chi^{(3)}I(\omega))}$$



• This yields for the refractive index of the medium:

$$n = \sqrt{\varepsilon_r} = \underbrace{\sqrt{(1 + \chi^{(1)} + \chi^{(3)}I(\omega))}}_{n_0}$$

• Which we can approximate as:

$$n = \sqrt{\varepsilon_r} \approx n_0 + \frac{1}{2n_0} \chi^{(3)} I(\omega)$$

• We set the *non-linear refractive index* of the material as $n_2 = \frac{1}{2n_0}\chi^{(3)}$, allowing us to express the, intensity dependent, refractive index as:

$$n(I) = n_0 + n_2 I(\omega)$$



Material	<i>n</i> ₀	$\chi^{(3)} (m^2/V^2)$	$n_2 ({\rm cm}^2/{\rm W})$	Comments and References b
Crystals				
Al ₂ O ₃	1.8	3.1×10^{-22}	2.9×10^{-16}	1
CdS	2.34	9.8×10^{-20}	5.1×10^{-14}	1,1.06 µm
Diamond	2.42	2.5×10^{-21}	1.3×10^{-15}	1
GaAs	3.47	1.4×10^{-18}	3.3×10^{-13}	1,1.06 µm
Ge	4.0	5.6×10^{-19}	9.9×10^{-14}	2, THG $ \chi^{(3)} $
LiF	1.4	6.2×10^{-23}	9.0×10^{-17}	1
Si	3.4	2.8×10^{-18}	2.7×10^{-14}	2, THG $ \chi^{(3)} $
TiO ₂	2.48	2.1×10^{-20}	9.4×10^{-15}	1
ZnSe	2.7	6.2×10^{-20}	3.0×10^{-14}	1,1.06 µm
Glasses				
Fused silica	1.47	2.5×10^{-22}	3.2×10^{-16}	1
As ₂ S ₃ glass	2.4	4.1×10^{-19}	2.0×10^{-13}	3
BK-7	1.52	2.8×10^{-22}	3.4×10^{-16}	1
BSC	1.51	5.0×10^{-22}	6.4×10^{-16}	1
Pb Bi gallate	2.3	2.2×10^{-20}	1.3×10^{-14}	4
SF-55	1.73	2.1×10^{-21}	2.0×10^{-15}	1
SF-59	1.953	4.3×10^{-21}	3.3×10^{-15}	1
Nanoparticles				
CdSSe in glass	1.5	1.4×10^{-20}	1.8×10^{-14}	3, nonres.
CS 3-68 glass	1.5	1.8×10^{-16}	2.3×10^{-10}	3, res.
Gold in glass	1.5	2.1×10^{-16}	2.6×10^{-10}	3, res.
Polymers				
Polydiacetylenes				
PTS		8.4×10^{-18}	3.0×10^{-12}	5, nonres.
PTS		-5.6×10^{-16}	-2.0×10^{-10}	6, res.
9BCMU			2.7×10^{-18}	7, $ n_2 $, res.
4BCMU	1.56	-1.3×10^{-19}	-1.5×10^{-13}	8, nonres, $\beta =$

TABLE 4.1.2 Third-order nonlinear optical coefficients of various materials ^a





Maria Goeppert Mayer predicts two photon absorption as part of her PhD – 1931 Theodore Maiman and the first laser (Ruby) – May 1960

	VOLUME 7, NUMBER 4 PHYSICAL REVIEW LETTERS AUGUST 15, 1961 GENERATION OF OPTICAL HARMONICS* P. A. Franken, A. E. Hill, C. W. Peters, and G. Weinreich The Harrison M. Randall Laboratory of Physics, The University of Michigan, Ann Arbor, Michigan (Received July 21, 1961) Second Harmonic Generation – 1961					
	Volume 7, Number	6 PHYSICAL REVIEW LETTERS SEPTEMBER 15, 1961				
1		TWO-PHOTON EXCITATION IN CaF ₂ :Eu ²⁺				
	W. Kaiser and C. G. B. Garrett Bell Telephone Laboratories, Murray Hill, New Jersey (Received August 28, 1961)					
nan and (Ruby)	Two photon absorption -1961					
	NONLINEAR OPTICS					
	4th Edition	Nicolaas Bloembergen, Nobel				
		Le de Discussion 1001				
		Laurate in Physics 1981, wrote				
	Nicolaas Bloembergen	NonLinear Optics in 1964				
	Sector Sector Sector					

World Scientific



FOUR-WAVE MIXING





 $\omega_3 \omega_1 \omega_2 \omega_4$

Image from Bryn A. Bell, Kai Wang, Alexander S. Solntsev, Dragomir N. Neshev, Andrey A. Sukhorukov, and Benjamin J. Eggleton, "Spectral photonic lattices with complex long-range coupling," Optica 4, 1433-1436 (2017) 53

 $\omega_3 \omega_1 \omega_2 \omega_4$

COHERENT ANTI-STOKES RAMAN SCATTERING



COHERENT ANTI-STOKES RAMAN SCATTERING









SOME PHYSICS...



BRAGG SCATTERING



evident that the peaks $A_1 B_1 C_1$, $A_2 B_2 C_2$ are analogous to spectra of the first and second orders, because of the absence of intervening sets of peaks. The value of n in the equation

$n\lambda = 2d\sin\theta$

seems clear. The difficulty of assigning a definite wave-length to the rays arises when we attempt to determine the value of d, the distance of plane from plane.

W. H. Bragg, W. L. Bragg, The Reflection of X-rays by Crystals, Proc. R. Soc. Lond. A 1913 88 428-438; DOI: 10.1098/rspa.1913.0040. Published 1 July 1913

The Nobel Prize in Physics 1915



Sir William Henry Bragg Prize share: 1/2 William Lawrence Bragg Prize share: 1/2

The Nobel Prize in Physics 1915 was awarded jointly to Sir William Henry Bragg and William Lawrence Bragg *"for their services in the analysis of crystal structure by means of X-rays"*

Photos: Copyright © The Nobel Foundation

*x-ray crystallography was also used to determine the double helix structure of the DNA molecule (1962 Nobel Prize in Physiology), haemoglobin (1962 Nobel Prize in Chemistry) and other proteins



BRILLOUIN SCATTERING



ELECTROSTRICTION



Force on a particle in an electric field:

$$F = -\nabla U = \frac{1}{2} \epsilon_0 \alpha_{eff} \nabla (E^2)$$
$$\alpha_{eff} = [(\alpha_{\parallel} - \alpha_{\perp}) \cos^2\theta + \alpha_{\perp}]$$



"The tendency of polarizable molecules to move towards regions of high optical electric field intensities is referred to as *electrostriction*."

R.W. Boyd, Nonlinear Optics. Elsevier, London, 3rd edition, 2008.

COHERENT RAYLEIGH-BRILLOUIN SCATTERING





Measurable quantities Temperature Density Speed of sound Shear & Bulk viscosity Polarizability <u>CRBS:</u>

- \checkmark Highly localized
- ✓ Laser beam signal output
- \checkmark Much lower resolvable densities



COHERENT RAYLEIGH-BRILLOUIN SCATTERING



Coherent Rayleigh Scattering

Jay H. Grinstead¹ and Peter F. Barker² ¹Optical Technology Division, National Institute of Standards and Technology, Gaithersburg, Maryland 20899 ²Department of Mechanical and Aerospace Engineering, Princeton University, Princeton, New Jersey 08544 (Received 21 January 2000)

VOLUME 89, NUMBER 18 PHYSICAL REVIEW LETTERS

28 OCTOBER 2002

Coherent Rayleigh-Brillouin Scattering

Xingguo Pan, Mikhail N. Shneider, and Richard B. Miles

Department of Mechanical and Aerospace Engineering, Princeton University, Princeton, New Jersey 08544 (Received 26 April 2002; published 10 October 2002)



A. Gerakis,¹ M. N. Shneider,² and P. F. Barker¹

¹Department of Physics and Astronomy, University College London, WC1E 6BT, United Kingdom ²Department of Mechanical and Aerospace Engineering, Princeton University, Princeton, New Jersey, 08544, USA

PHYSICAL REVIEW A 87, 033825 (2013)

Narrowband coherent Rayleigh-Brillouin scattering from gases confined by a high-intensity optical lattice

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COHERENT RAYLEIGH-BRILLOUIN SCATTERING Single Shot

<u>Chirped</u> <u>lattice</u>

> CRBS: ✓ Highly localized ✓ Laser beam signal output ✓ Much lower resolvable densities ✓ Single Shot LUXEMBOURG INSTITUTE OF SCIENCE AND TECHNOLOGY

Gerakis, A., Shneider, M.N. and Barker, P.F., 2013. Singleshot coherent Rayleigh–Brillouin scattering using a chirped optical lattice. Optics letters, 38(21), pp.4449-4452.

COHERENT RAYLEIGH-BRILLOUIN SCATTERING Single Shot





COHERENT RAYLEIGH-BRILLOUIN SCATTERING Experimental Setup



COHERENT RAYLEIGH-BRILLOUIN SCATTERING Experimental Setup



Dual color, frequency, pulse duration and shape agile laser system for particle spectroscopy and manipulation

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Display 🖉 Cursors 🗉 Measure 🖬 Math 🗠 Analysis 🛠 Utilities 🐽 Sund



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COHERENT RAYLEIGH-BRILLOUIN SCATTERING Experimental Setup

Experimental parameters

>450 mJ/pulse/arm
> Up to ~10 GHz chirp rate
> Variable pulse durations: 10-1000ns
> Rep. Rate: 10-30 Hz







Arc method :

- Simple to implement
- High nanomaterial yield
- Variety of synthesized

nanostructures

Environmental conditions :

Helium pressure	-500 Torr
Electrodes	Carbon, <1 cm in diameter
Power	few kW
Plasma density	$10^{14} - 10^{16} \mathrm{~cm}^{-3}$
Temperature	~1 eV
Plasma role in na	nostructure synthesis?

- What are plasma properties?
- How feedstock material is formed?
- What growth conditions are realized in the arc?

Cathode Anode

Particles are formed here

Imaging through C2 filter

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Goal: create prescribed nanoparticles

- Model arc synthesis: plasma modeling and growth modeling
- <u>In situ</u> diagnostics to validate model
 - No optical way to measure nanoparticles from <u>initial nucleation to</u> <u>~10nm</u>... unless we use <u>coherent Rayleigh-Brillouin scattering</u>





APPLIED PHYSICS LETTERS 102, 173109 (2013)

CrossMark

Application of coherent Rayleigh-Brillouin scattering for *in situ* nanoparticle and large molecule detection

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- The arc run time is ~ 2 minutes
- Nanoparticles have got a residence time of ~1 ms [1]
 - Residence time of nanopartilces affects growth



[1] M. Kundrapu, M. Keidar, Numerical simulation of carbon arc discharge for nanoparticle synthesis, Physics of Plasma, Vol. 19, No. 073510, 2012







A. Gerakis, Y-W. Yeh, M. N. Shneider, J. M. Mitrani, B. C. Stratton, and Y. Raitses, "Four-Wave-Mixing Approach to *In Situ* Detection of Nanoparticles", Phys. Rev. Applied 9, 014031 (2018)

 $\rm C_{60}$ FWHM (@1500 K): 309.9 m/s Estimated FWHM (@1500 K): 18.8 m/s

- Detected nanoparticles of: ~500.000 AMU
- C_{60} : 720 AMU

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- Assuming the same density, the detected nanoparticles' dimensions are: ~6 nm!!!
- Assuming polarizability the same as C_{60} then measured density is: ~10¹² cm⁻³

For the first time detection of nanoparticles of dimensions of <10 nm, <u>in situ</u> in plasma.

• Now we are able to map out in situ nanoparticle nucleation and growth in the arc discharge, with high spatial resolution (~100 μ m).

A. Gerakis, Y-W. Yeh, M. N. Shneider, J. M. Mitrani, B. C. Stratton, and Y. Raitses, "Four-Wave-Mixing Approach to *In Situ* Detection of Nanoparticles", Phys. Rev. Applied 9, 014031 (2018)


COHERENT RAYLEIGH-BRILLOUIN SCATTERING Neutral gas flow measurement



COHERENT RAYLEIGH-BRILLOUIN SCATTERING Neutral gas flow measurement



COHERENT RAYLEIGH-BRILLOUIN SCATTERING Neutral gas flow measurement







COHERENT RAYLEIGH-BRILLOUIN SCATTERING Neutral gas flow measurement Compressed





Alexandros Gerakis, Junhwi Bak, Robert Randolph and Mikhail N. Shneider. "Demonstration of single shot laser velocimetry with coherent Rayleigh-Brillouin scattering," AIAA 2021-0224. *AIAA Scitech 2021 Forum.* January 2021.



COHERENT RAYLEIGH-BRILLOUIN SCATTERING Two color CRBS scheme



♦ More efficient Rayleigh scattering $(\propto \lambda^{-4})$

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- Measure simultaneously neutral density and translational temperature in glow discharge
 - Two color, single shot
 CRBS as means of
 measurement

 - Scan radially across the discharge.









Dual color CRBS chamber setup

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- Measurements radially across the positive column of glow discharge, from center to periphery.
 - Assumption: Positive column is radially symmetric
 - Assumption: Operating conditions are the same from point to point and run to run.
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Xenon 15 Torr @300K

- Temperature: $f(v_z) \propto e^{\frac{-mv_z^2}{2kT}}$
- Density: $I_S \propto \Delta n^2$







Pan, Xingguo & Shneider, M.N. & Miles, R.B. (2004). Coherent Rayleigh-Brillouin scattering in molecular gases. Physical Review A. 69. 33814-. 10.1103/PhysRevA.69.033814.





COHERENT THOMSON SCATTERING A novel way to measure electron temperature?



Physics of Plasmas

ARTICLE scitation.org/journal/php

Analysis of coherent Thomson scattering from a low temperature plasma

Cite as: Phys. Plasmas **29**, 033507 (2022); doi: 10.1063/5.0072540 Submitted: 23 September 2021 · Accepted: 21 February 2022 · Published Online: 11 March 2022





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LASER DIAGNOSTICS Summary

- ➤ There's no "one-size fits all": different techniques are complimentary towards the solution of the same problem.
- Ideal way to characterize plasma & plasma flow measurements.
- \succ Work with (almost) every gas.
 - \succ Advantage <u>and</u> disadvantage
- \succ Time to replace all mechanical probes $\textcircled{\sc {\odot}}$







Thank you! Questions?

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