

## Plasma turbulence

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computer simulation of magnetized plasma turbulence



## About me:



## Topics in this talk:

- Turbulence: what it isn't, what it is. & examples.
- Instability as a cause of turbulence: Rayleigh-Taylor as an example.
- Turbulent energy cascades.
- Examples of turbulence in plasmas.
- Mechanisms that can drive turbulence in plasmas.
- Ways to explore turbulence: experiment & computer simulation (including ML)



### **Turbulence**: Everywhere in nature & in man-made fluids. We all experience it.

Let's start by seeing its onset in a channel with fluid past an obstacle oozing green dye:



Credit: <u>The Lutetium Project</u>

Visually we can say turbulence is:

- Complex
- Hard to predict
- Swirly



### **Turbulence:** Everywhere in nature & in man-made fluids. We all experience it.



### **Turbulence:** Everywhere in nature & in man-made fluids. We all experience it. Yet, how do we define it?

Cool at-home demo:

visualize air turbulence in your room using •

- A fog machine.
- Diffracted laser light.



photons scatter off of air molecules and reach our eyes, allowing us to see the turbulent structures in the air

Credit: <u>3blue1brown</u>

Other properties of turbulence more closely related to its mathematical description:

- Chaotic.
- Ergodic.

- Diffusive and mixing.
- Unstable



### Instability is key: Turbulence arises in systems in unstable equilibrium, where perturbations can grow and lead to instability



Recall the position of the pendulum:  $x(t) \propto Ae^{i\omega t}$ 

Fluid mechanics turbulence is often studied via the Navier-Stokes equation. For incompressible fluids:

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$$\begin{array}{l}
\rho \frac{D\mathbf{u}}{Dt} = \rho \left( \frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla \right) \mathbf{v} = -\nabla p + \nu \rho \nabla^2 \mathbf{u} + \rho \mathbf{g} \\ \text{inertia} \\ \text{(like } ma \text{ in } F = ma) \end{array} \stackrel{\text{pressure}}{\text{force}} viscous \\ \text{force} \quad \text{for$$

### Example of an unstable fluid: dense fluid on top, light fluid on the bottom



done with **Gkeyll**,

for modeling plasmas

a free, open-source code

How does density stratification lead to instability?

$$p\left(\frac{\partial}{\partial t} + \mathbf{v} \cdot \nabla\right) \mathbf{v} = -\nabla p + \nu \rho \nabla^2 \mathbf{u} + \rho \mathbf{g}$$
$$\frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T = 0 \qquad \nabla \cdot \mathbf{v} = 0$$

Neglect viscosity and nonlinearities, one can combine these equations to show that the frequency of sinusoidal perturbations (waves) in this scenario is given by

$$\omega = \sqrt{-\frac{g}{\rho}\frac{\partial\rho}{\partial z}}$$

and if we assume  $\rho = e^{y/L}$ 

$$\omega = i \sqrt{\frac{g}{L}} \quad \Rightarrow \quad \gamma = \sqrt{\frac{g}{L}}$$

Compare with pendulum in previous slide.





### The perturbations grow until fully developed turbulence ensues



Credit: Mark Stock (Vic2d)



Many other sources of instability in fluids; e.g. the Kelvin-Helmholtz is driven by gradients in the fluid velocity.



# As turbulence evolves, we see the formation of smaller eddies, and a rich distribution of eddies of all sizes.

In the home-demo visualization of air turbulence in your room:



Credit: <u>3blue1brown</u>



# As turbulence evolves, we see the formation of smaller eddies, and a rich distribution of eddies of all sizes.

The "swirly" nature of fully developed turbulence is due to a transfer of energy from large scales structures in the fluid, to smaller and smaller swirls (eddies)



known as an "energy cascade".

This continues to ever-smaller scales, until molecular dissipation converts kinetic energy of the eddies into thermal energy (heating up the fluid). The energy flow in a turbulence fluid is



# Determination of the rate of energy transfer between scales is one of the most consistent results in turbulence research

At some scales smaller than the large scale of the external forces (or driving flow inhomogeneities), but larger than the dissipative scale, this transfer of energy must be constant. Otherwise we'd see a build up of kinetic energy at one particular scale.

Russian scientist Andrey Kolmogorov determined the spectrum of the energy as a function of the size of the eddies. Similar to how special relativity stems from three basic assumptions (c=constant, c is the max speed, inertial universality of the laws of physics), Kolmogorov employed two basic assumptions<sup>1,2.</sup>

He was concerned with differences in velocity between two points

 $\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t) = \mathbf{u}(\mathbf{x} + \mathbf{r}, t) - \mathbf{u}(\mathbf{x}, t)$ 

in stationary, homogeneous and *isotropic* turbulence.





<sup>1</sup> Kolmogorov, A. N. Dokl. Akad. Nauk SSSR, 32, 19-21 (1941).
 <sup>2</sup> Lévêque, E. EAS Publications Series 21, 7 (2006).
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The two hypotheses were:

1. At  $r \ll L$  the distributions of  $\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t)$  are "universal", i.e. independent of the type of external force or driving mechanism, & only a function of  $\nu$  (viscosity) and the mean energy dissipation rate per unit mass ( $\overline{\epsilon}$ ).

Units 
$$\begin{bmatrix} \nu \end{bmatrix} = m^2/s$$
  
 $[\overline{\epsilon}] = J/(kg s) = m^2/s^3$   $\left\{ |\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t)|^2 \right\} = \sqrt{\nu \overline{\epsilon}} \Phi(r/\ell_d)$ 

can also infer the dissipation length scale  $\ \ell_d = (
u^3/ar{\epsilon})^{1/4}$ 

2. At  $\ell_d \ll r \ll L$  the distributions of  $\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t)$  are independent of  $\nu$ .

$$\left< |\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t)|^2 \right> = b \left( \overline{\epsilon} \, r \right)^{2/3} = b \, \overline{\epsilon}^{2/3} \, k^{-2/3} \qquad \qquad k: \text{wavenumber}$$

From this last equation deduce the energy spectrum

 $\mathcal{E}(k) = C \,\overline{\epsilon}^{2/3} \, k^{-5/3}$ Kolmogorov spectrum

since

$$\left\langle |\delta \mathbf{u}(\mathbf{x}, \mathbf{r}, t)|^2 \right\rangle \propto E = \int_0^\infty \mathcal{E}(k) \, \mathrm{d}k$$

<sup>1</sup> Kolmogorov, A. N. Dokl. Akad. Nauk SSSR, 32, 19-21 (1941). <sup>2</sup> Lévêque, E. EAS Publications Series 21, 7 (2006).

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average

function

dimensionless

 $\Phi$ : Universal

## Studying, modeling & explaining turbulence is hard. But Kolmogorov's spectrum is a successful result with countless experimental & numerical confirmations

#### $10^{7}$ $10^{6}$ $(k/k_{\rm d})^{-5/3}$ $10^{5}$ $10^{4}$ $10^{3}$ $\mathcal{E}(k)$ $10^{2}$ 10 $R_{\lambda}$ Physical situation 1 2000 Tidal channel $10^{-1}$ Boundary layer 850 Wake behind cylinder Pao's (1965) 308 Pipe flow 170 correlation $10^{-2}$ Homogeneous shear flow -130Grid turbulence 72 Wake behind cylinder 23 $10^{-3}$ Grid turbulence $10^{-}$ $10^{-2}$ $10^{-5}$ $10^{-3}$ $10^{-1}$ 10 $10^{-4}$ $k \ell_d$

Navier-Stokes & Euler simulations of a dense gas<sup>2</sup>:





<sup>1</sup> McComb, W. D. The Physics of Fluid Turbulence. Oxford University Press (1990). <sup>2</sup> Sciacovelli, L., et al. J. Fluid Mech. 800, 140 (2016).

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Experiments<sup>1</sup>:

# Kolmogorov's spectrum has even been accounted for in non-scientific matters...

### Van Gogh's La Nuit étoilée



"the probability distribution function (PDF) of luminance fluctuations in some impassioned van Gogh paintings, painted at times close to periods of prolonged psychotic agitation of this artist, compares notable well with the PDF of the velocity differences in a turbulent flow as predicted by the statistical theory of Kolmogorov"<sup>1</sup>

<sup>1</sup> J. L. Aragón, et al. J. of Mathematical Imaging & Vision 30, 275 (2008).
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Turbulent energy transfer is usually from large to small scales, but there are instances of "inverse energy cascades"



Jupiter's turbulent atmosphere forms organized "jet streams"

Somehow energy is transferred from small scales to these large scale flows.

#### F: 115 T: 1.1500e+02



### **Plasma turbulence is everywhere laboratories & nature**



Ä

### Other similarities exist between neutral fluid & plasma turbulence? Instability drives turbulence. Power law energy spectrum.

There are many drivers of instability in plasmas. Let's consider 1 from tokamaks:

denser plasma magnetic field (B) less dense vro-orbits of plasma charged particle Plasma feels an effective gravitational acceleration due to the centrifugal force:  $g_{\text{eff}} = -\frac{1}{R}$ 

Toroidal plasma (bird's-eye view)

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This scenario has all the ingredients of the Rayleigh-Taylor instability we studied above:

- (effective) gravity
- denser "fluid" on top of lighter "fluid"

So recycling our earlier result, we'd say that the growth rate is

$$\gamma = \sqrt{\frac{g_{\text{eff}}}{L}} = \frac{v}{\sqrt{RL}}$$

The length scale *L* here turns out to be related to the pressure gradient, e.g.  $L = p/|\nabla p|$ 

Is the mechanism the same as in a neutral fluid?



Consider a patch of plasma on the outboard region of the torus:

Particles will drift with velocity:





Consider a patch of plasma on the outboard region of the torus:





Consider a patch of plasma on the outboard region of the torus:



This  $E \times B$  drift reinforces the perturbation, creating an (unstable) positive feedback.





- These perturbations (eddies) have a size of order  $\sim r_L = m v_{\perp} / (|q|B).$
- $\mathbf{v}_{\mathbf{E}} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} \quad \text{o Particles/eddies move} \sim r_L, \\ \text{causing } n \& T \text{ to redistribute the} \\ \text{eddy to vanish.}$ 
  - $\Rightarrow$  eddy lifetime

 $au_{\mathrm{eddy}} \sim r_L / v_{\mathbf{E}}$ 

- Can obtain a diffusion time estimate. Consider a diffusion equation  $\frac{\partial f}{\partial t} = D \frac{d^2 f}{dx^2}$ , the diffusive time scale is  $\tau_d \sim L^2/D$ . So we can say that energy diffusion due to the above process takes place on the timescale:  $\tau_d \sim \frac{L^2}{r_L^2/\tau_{eddy}} \sim \frac{L^2}{r_L \mathbf{v_E}}$   $[D] = m^2/s$
- How large is  $\mathbf{v_E}$ ?  $\mathbf{v_E} \sim \frac{1}{BL} \phi \sim \frac{1}{|q|BL} T = \frac{v_t^2}{L\omega_c} = \frac{r_L}{L} v_t$  $\Rightarrow \tau_d \sim \frac{L^2}{r_L^2} \frac{L}{v_t}$  estimate for confinement time  $\tau_E = \frac{\text{plasma energy}}{\text{rate of energy loss}}$

# Energy is being lost from the fusion-producing plasma core. Is all lost? How do we stabilize RT-like & micro instabilities?



A different behavior of between good & bad curvature regions is clearly observed in computer simulations of tokamak turbulence



# Energy is being lost from the fusion-producing plasma core. Is all lost? How do we stabilize RT-like & micro instabilities?

Turbulent transport can be mitigated by sheared flows:

F: 120 T: 1.2000e+02



Computer simulation done with **Gkeyll**, <u>a free, open-source code for modeling plasmas</u>



# Studying, understanding and predicting turbulence is important & can have a significant impact on the cost of fusion

As an example, take a study of the sensitivity of a compact fusion power plant's capital cost:



<sup>1</sup>M. R. Wade, J. A. Leuer. Fus. Sci. And Tech. 77, 119 (2021).



### **Turbulence is frequently studied via numerical simulation**

Sophisticated computer codes that solve fluid, kinetic or gyrokinetic equations (or all of the above as is the case of the **Gkeyll** code I work on) have been developed around the world.

e.g. a gyrokinetic simulation of electromagnetic plasma turbulence with **Gkeyll**<sup>1</sup>:

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There are numerous other codes that can be used to study plasma turbulence, e.g.:

- GENE, GYRO, GS2, GX, Stella, GKV, AstroGK, XGC.
- GRILLIX, TOKAM3X, GBS.
- Athena
- Vlasiator
- Many, many others



# Plasma turbulence is also explored experimentally. For example, via gas puff imaging

Electrons collide with puffed neutral gas, leading to ionization events and emitted light.





These collisions are more frequent where the local density is higher, allowing us to visualize plasma fluctuations

<sup>1</sup>S. J. Zweben, et. al. Review of Sci. Inst 88 (2017).
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# Increasingly, turbulence is also being studied w/ machine learning

One out of many examples of machine learning applied to studying turbulence is the development of physics-informed neural networks (PINN)<sup>1</sup>:



Authors trained a PINN using data from these locations with a loss function that incorporates the incompressible Navier-Stokes equation. Then used the PINN to advance the solution (velocity field), and obtained reasonable agreement with the Navier-Stokes-only simulations.



<sup>1</sup>M. Raissi et al. JCP 378, 686 (2019).

# Summary

- Fluid turbulence is ubiquitous, and often driven by a system in unstable equilibrium.
- Fluid turbulence exhibits a range of scales, with an energy cascade between them that is characterized by the Kolmogorov spectrum:

 $\mathcal{E}(k) = C \,\overline{\epsilon}^{2/3} \, k^{-5/3}$ 

- Some plasma instabilities have close analogs in plasmas, as is the case of the Rayleigh-Taylor instability.
- Instabilities lead to fully developed turbulence & a power law energy spectrum in plasmas too, but the turbulent transport can be mitigated by clever design of magnetic fields and sheared flows.
- Sophisticated computer codes exist for studying turbulent plasmas.
- Turbulence can strongly impact the cost of a fusion power plant.

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a free, open-source code for modeling plasmas

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## Additional references:

Neutral fluid turbulence:

- U. Frisch, Turbulence: The legacy of A.N. Kolmogorov (1995).
- S. Pope, *Turbulent flows* (2000).

Plasma turbulence:

- P. H. Diamond, *Modern plasma physics: Volume 1, Physical Kinetics of Turbulence Plasmas* (2010).
- R. J. Goldston & P. H. Rutherford, Introduction to Plasma Physics (1995).
- J. Krommes, *The gyrokinetic description of microturbulence in plasmas.* Annual Review of Fluid Mechanics 44, 175 (2012).

Also, if you are interested in data-mining turbulence simulations, check out the <u>John</u> <u>Hopkins Turbulence Database</u>.

