Introduction to the physics of complex/dusty plasmas

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Outline

- Why dusty plasmas?
- Basic properties
 - Fundamental parameters
 - Charging
 - Forces and transport
- Dusty plasma studies
 - Statistical mechanics of non-equilibrium systems
 - Dust acoustic / dust density waves
 - Magnetic field effects
- Outlook
 - Outstanding problems and opportunities
- More information
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WHY DUSTY PLASMAS?





Dusty (complex, fine particle, colloidal) plasmas

- Complex plasmas four component plasma system
 - Ions
 - Electrons
 - Neutral atoms
 - Charged microparticles
- Plasma and charged microparticles coupled via collection of ions and electrons from the background plasma.
- Presence of microparticles:
 - Modifies density and charge distribution
 - Modifies plasma instabilities
 - Introduces new dust-driven waves
 - Ubiquitous in natural and man-made plasmas
- Measurements of dust particles:
 - Forces
 - Electrostatic potential
 - Velocity distributions







Dusty plasmas in astrophysical environments



http://hubblesite.org/newscenter/archive/ releases/2007/16/image/f/format/large_web/

Image: Star formation in Carina Nebula



http://www.almaobservatory.org/en/pressroom/press-releases/771-revolutionaryalma-image-reveals-planetary-genesis

Image: HL Tau (2014)

Photoionization from stellar material charges the dust in the nebula. The presence of charged dust may lead to enhanced coagulation of small particles AND to repulsion between larger particles. [F. Verheest, PPCF, **41**, A445 (1999)]





Dusty plasmas in the solar system



- Discovered by Voyager 2 in 1980
- Spokes seen in forward scattered light –> composed of micron-sized dust and ice
- Spokes exhibit dynamical behavior on timescales of minutes.





Dusty plasmas in the solar system

- Cassini (2004) initially, no spokes observed!
- Re-evaluation of the Voyager data to understand changes in plasma environment.
- Cassini (2006) reappeared as predicted by updated models.
- Spoke formation correlated to opening angle of Saturn's rings.



 $1980, B = 6^{\circ}$





Cassini observation of spokes C. J. Mitchell, et al., Science, **311**, 1589 (2006)



Aug 21, 2008

Dusty plasmas in terrestrial environments

- Noctilucent clouds (NLC's) form at extremely high altitudes, about 85 km, that "shine at night".
- They form in the cold, summer polar mesopause and are believed to be charged ice crystals.
- They are believed to be associated with radar backscatter phenomena (PSME's) observed during the northern summers.







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Dust contamination:

- During the 1990s research was driven by the formation of microparticles in plasma processing reactors.
- Here, microparticles up to several microns in diameter can be grown in the plasma.
- "Killer" particle size has diam. \leq 20 nm.







AUPSEOM: http://fjwsys.lanl.gov/bpw/contamination.html - G. Selwyn, LANL



Dust in fusion devices

- Microparticles (d ~ 10 nm) entering the main plasma of a toroidal device can degrade plasma performance and possibly induce disruptions.
- Dust production of up to several kg's / day could occur in a large scale, ITER-like fusion device.
- Photographs show micron-sized dust particles formed in TEXTOR-94.



0.1 mm





J.Winter, PPCF, **40**, 1201 (1998) S. I. Krasheninnikov, et al., PPCF, **53**, 083001 (2011)



Dust in fusion devices

from \EFIT01, Shot 111877, time= 67.00ms



- Dust in fusion plasmas was often not considered a major issue.
- Recent studies do indicate that uncontrolled dust injection can have detrimental impacts on fusion plasmas.





Dust contamination in Iter



- Tritium retention in "dust"
- Health and safety hazard
- Reduction of density control
- Degradation of first wall material
- Radiated power losses



BASIC PROPERTIES





Basic properties of dusty plasmas

- Fundamental parameters
- Charging
- Forces





What are the parameters of a dusty plasma?





Fundamental Parameters (1): The basic equations

- Define: Relevant scales for a dusty plasma
- Use continuity and momentum equations
 - Assume no zero-order gradients or flows
 - Assume only electrostatic oscillations
 - Close set of equations using Poisson's equation.

Parameters:

$$s - ion, elec, dust$$

 $a - dust radius$
 $q_s - charge; q_d = -Z_de$
 $n_s - density$
 $m_s - mass$
 $v_s - velocity$
 ϕ - potential

- Continuity:
- Momentum:
- Poisson's:

$$\frac{\partial n_s}{\partial t} + \nabla \cdot \left(n_s \vec{v}_s \right) = 0$$

$$\frac{\partial \vec{v}_s}{\partial t} + \left(\vec{v}_s \cdot \nabla \right) \vec{v}_s = -\frac{q_s}{m_s} \nabla \phi$$

$$\nabla^2 \phi = -\frac{1}{2} \sum q_s n_s$$

ε



Fundamental Parameters (2): Large mass extends the time scales

- Linearize the equations using: $a = a_0 + a_1 \exp[i(kx-\omega t)]$
- Derive a result that gives the time scales of plasma oscillations:

$$\omega^{2} = \sum \omega_{ps}^{2} = \omega_{pe}^{2} + \omega_{pi}^{2} + \omega_{pd}^{2}$$
where:
$$\omega_{ps}^{2} = \frac{q_{s}^{2}n_{0s}}{\epsilon_{0}m_{s}}$$

• For typical lab plasma parameters: $f_{ps} = \omega_{ps}/2\pi$ $n_{i0} = n_{e0} \sim 10^{14} \text{ m}^{-3}$, $n_{d0} \sim 10^{10} \text{ m}^{-3}$, argon plasma, $Z_d \sim 4600$, $a \sim 1.5 \mu \text{m}$

•
$$f_{pe} = 90 \text{ MHz}$$
, $f_{pi} = 330 \text{ kHz}$, $f_{pd} = 23 \text{ Hz}$





Fundamental Parameters (3) – Coupling parameter is a measure of self-organization

- Γ (coupling parameter) is indicative of the selforganizing, emergent properties of dusty plasmas.
- A dusty plasma can be used as a model system to investigate problems in soft-matter physics.
- Assume dust particles interact via a screened Coulomb interaction (Yukawa, Debye-Hückel): $\varphi \sim \frac{\exp(-r / \lambda_D)}{r}$

$$\Gamma = \frac{\text{electrostatic potential energy}}{\text{thermal energy}} = \frac{Q_d^2}{4\pi\varepsilon_0 kT_d\Delta}$$

$$\Delta = \text{Wigner-Seitz radius} = \left(\frac{4\pi n_d}{3}\right)^{-1/3}$$





Redefining the parameters of a dusty plasma







G. Morfill, et al., PoP, 6, 1769 (1999).

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Basic properties of dusty plasmas

- Fundamental parameters
- Charging
- Forces





Fundamental Parameters (4) – Dust grain charge is a dynamic variable





- A dynamic equilibrium is established as the grain electrically floats in the plasma: $I_{total} = I_{electron} + I_{ion} + I_{see} + I_{thermionic} + I_{hv} = f(n_j, T_j, \phi; \underline{r}, t)$
- Implication: dQ_d/dt ≠ constant;
- Grain charge ($Q_d = Z_d e$) is a new dynamic variable





Estimating dust grain charge (1)

- For laboratory studies, ions and electrons are the dominant charging mechanisms.
- We assume dust behaves as an electrically floating probe and estimate the flux to the grain surface using orbit motion limited (OML) theory.

$$I_e = 4 \pi a^2 \left(\frac{e n_e}{4}\right) \left(\frac{8 k T_e}{\pi m_e}\right)^{1/2} \exp\!\left(\frac{e U}{k T_e}\right) \quad \text{electron}$$

- a grain radius
- T_s temperature
- k Boltzmann's constant
- U grain surface potential

$$I_i = 4 \pi a^2 \left(\frac{e n_i}{4}\right) \left(\frac{8 k T_i}{\pi m_i}\right)^{1/2} \left(1 - \frac{e U}{k T_i}\right)$$
 ion



Estimating dust grain charge (2)

- Assume grains are conducting.
- Assume grains are spherical capacitors: $Q_d = \pm eZ_d = 4\pi\epsilon_0 aU$
- Assume quasineutrality: $en_i = en_e + Q_d n_d$
- Solve the balance equation: $I_e + I_i = 0$

$$\left(1+\frac{Q_{d}n_{d}}{en_{0}}\right)\left(\frac{m_{i}T_{e}}{m_{e}T_{i}}\right)^{\frac{1}{2}}\exp\left(\frac{eU}{kT_{e}}\right)=1-\left(\frac{eU}{kT_{i}}\right)$$

• Solve numerically for the grain potential U to get the charge, Q_d .





Charging experiments (1)



Charging experiments (2)

Barkan experiment uses a Q-machine to generate the plasma -> Here, $T_i \approx T_e \sim 0.2 \text{ eV}$



FIG. 2. Schematic diagram of the device used to disperse dust into the plasma column.



FIG. 3. Langmuir probe characteristics obtained under identical conditions, except for the absence (upper plot) or presence (lower plot) of kaolin dust. In the lower characteristic, the dust dispenser is abruptly turned off near the end of the trace to check that the electron current returns to the no-dust value.

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A. Barkan, et. al., Phys. Rev. Lett., 73, 3093 (1994)

Charging experiments (3)



Grain charging in a dc glow discharge plasma



FIG. 3. Electron saturation current measurement as a function of axial position in the FPS device. The open circles indicate measurements of the electron saturation current in the absence of the silica dust and the closed squares indicate measurements of the electron saturation current in the presence of the silica dust particles.

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E. Thomas and M. Watson, Phys. Plasmas, 7, 3194 (2002)

Basic properties of dusty plasmas

- Fundamental parameters
- Charging
- Forces





Summary of dominant forces in dusty plasmas

Force	Origin	Size dependence
Weight	Gravity	a ³
Neutral drag	Streaming neutrals	a ²
lon drag	Streaming ions	a ²
Thermophoretic	Temperature gradient	a ²
Electric	Electric field	a'
Magnetic	Magnetic field	a'

These forces give rise to the majority of the phenomena observed in laboratory and microgravity dusty plasma experiments a = dust grain radius

Adapted from textbook Plasma Physics by A. Piel, Table 10.2 (Springer-Verlag, 2010)





Gravitational force

 Relevant for dusty / complex plasmas since the particles need to be suspended in the plasma.

•
$$\vec{F}_{\rm g} = m_{\rm d}\vec{g} = \frac{4}{3}\pi a^3\rho_{\rm d}\vec{g}$$

• Gravitational force can also play an important role in astrophysical environments (e.g., Saturn's rings).







Electric force

 Because the dust grains are charged, they respond to the internal electric fields within the plasma.

•
$$\vec{F}_{\rm E} = Q_{\rm d}\vec{E} = 4\pi\epsilon_0 a\phi_{\rm fl}\vec{E}$$

 Ion flows can cause the shielding cloud around dust grains to become distorted, leading to a dipole-like charge distributions.

 $\vec{F}_{\rm dip} = \vec{\nabla} (\vec{p} \cdot \vec{E})$





Ion flows can lead to a downstream positive wake below a dust grain



Balancing gravitational and electric forces: A zero-order equilibrium in laboratory dusty plasmas

- For most ground-based experiments: F_{gravity} ~ F_{electric}
- Defines zero-order equilibrium
- Typical values:
 - → a = 1.5 µm
 - \Rightarrow m_d = 2.8 x 10⁻¹⁴ kg
 - \Rightarrow Z_d = 4600 electrons
 - → E = 3.8 V/cm





lon drag force - origin

- Positive ions can flow in a plasma in the direction of the electric field.
- Arises from the momentum transfer from ion-dust interactions: the collection of ions and Coulomb collisions.
- $F_{ion-drag} = F_{collection} + F_{collision}$
- Critically depends upon the screening length of the dust particle.



Refs: M. Barnes, et al., PRL, 68, 313 (1992)
S. Khrapak, et al., PRE, 66, 046414 (2002)
A. Ivlev, et al., PRE, 71, 016405 (2004)
I. Hutchinson, et al., PPCF, 48 185 (2006)
S. Khrapak, et al., IEEE TPS, 37, 487 (2009)



Evidence for ion drag force: voids in space

- In rf plasmas in microgravity experiments, a void appears.
- Enhanced ionization at the center: provides a source of ions and outward electric field.
- Boundary determined by the force balance: $F_{elec} = F_{ion-drag}$







Evidence for ion drag force: voids in the lab

- Control size of the void region using different potentials on a probe tip.
- Estimate the void size, x₀, using different electric field estimates.

Dashed: 25 V/cm Solid: 17 V/cm





Thermophoretic force

- Arises from the momentum transfer from neutral-dust interactions.
- Neutral atoms from "hot" side provide more momentum than those from "cold" side.

$$F_{thermo} = -\frac{8}{3} \frac{a^2}{v_{tn}} \Lambda \frac{dT_n}{dz}$$

 v_{tn} – neutral thermal velocity T_n – neutral gas temperature Λ – thermal conductivity







Evidence for thermophoretic force: simulated voids



- Dusty plasma of 3.4 µm melamine formaldehyde particles
- Applied temperature gradient of ~1200 K/m
- $\Delta T = 25 \text{ K over } 20 \text{ mm}$
- Lower electrode is heated using a Peltier element



H. Rothermel, et al., PRL, 89, 175001 (2002)
Evidence for thermophoretic force: Coulomb balls



O.Arp, et al, Phys. Plasmas, 12, 122102 (2005)

Combinations of these forces give rise to transport phenomena in dusty plasmas

[Auburn University - Plasma Sciences Laboratory]

DUSTY PLASMA STUDIES: STATISTICAL MECHANICS

A distribution function can represent the state of a system as it involves in time and space

- Assume that a measurement can be made of the position and velocity of each atom at time t₀.
- The distribution function represents the probability of finding a particle with position vector (<u>r</u>) and velocity vector (<u>v</u>).
- Distribution function: $f(\underline{v}, \underline{r}, t)$

A distribution function can represent the state of a system as it involves in time and space

- Integrating the distribution function (moments) can give macroscopic quantities:
- Number: $n(\vec{r},t) = \int f(\vec{r},\vec{v},t) d^3 \vec{v}$
- Momentum: $n\vec{p} = \int m\vec{v}f(\vec{r},\vec{v},t)d^3\vec{v}$
- Kinetic energy: $\frac{1}{2}m\langle \vec{v}^2 \rangle = \int \frac{1}{2}m\vec{v}^2 f(\vec{r},\vec{v},t)d^3\vec{v}$
- Time evolution:

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + v \bullet \nabla f + \left(\frac{\vec{F}}{m}\right) \bullet \nabla_v f = \left(\frac{\partial f}{\partial t}\right)_{collisions}$$

In a dusty plasma it is possible to directly measure the complete distribution function - $f(\underline{r}, \underline{v}, t)$

In a dusty plasma it is possible to directly measure the complete distribution function - $f(\underline{r}, \underline{v}, t)$

Studies of the thermodynamics of dusty plasmas

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Studies of statistical properties of dusty plasmas

Experiment	Reference
Melting of plasma crystals	Melzer, et al., PRE (1996)
Melting of plasma crystals	Quinn, et al., PoP (2000)
Dusty plasma liquids	YJ. Lai et al, PRL (2002)
Microstructures in 2d dusty plasma liquids	YH. Huang, et al., PRE (2007)
Propagation of melting fronts	Williams, et al., PRE (2012)
Laser driven heating of plasma crystals	Nosenko, PoP (2006)
Superdiffusion	Liu, PRE (2008)
Viscoelasticity in 2-d liquids	Feng, PRL (2010)
Thermal properties of weakly-coupled system	Zhakhovskii, JETP Lett. (1997)
Thermal properties of weakly-coupled system	Williams, et. al., PoP (2007)
Thermal properties of weakly-coupled system	Thomas, et al., PoP (2007)
Anisotropic velocity distributions	Fisher, et al., PRE (2012)

 Since the coupling parameter, Γ, can be controlled, a dusty plasma is a model system for studying the statistical mechanics of solids, liquids, gases and plasmas.

rf plasma laser excited dc plasma

DUSTY PLASMA STUDIES: DUSTY PLASMA WAVES

Waves in dusty plasmas

- In dusty plasmas, it is possible to visualize both collective motion and individual particle motion simultaneously.
- This enables a unique capability to study wave properties with incredible precision.
- Dusty plasma waves can be studied from "strongly-coupled" (solid-like) systems to "weaklycoupled" (fluid-like) systems.

Dust wave in a strongly coupled dust lattice

Laser excited one-dimensional and two-dimensional lattice waves in a plasma crystal.

Dust density waves in laboratory and microgravity plasmas

Dust density waves refers to the general class of low frequency, often selfexcited waves in a dusty plasma that are characterized by a modulation of the dust number density.

Left: DC glow discharge experiment (Auburn) Above: RF microgravity experiment (Kiel)

Existence of a "dust sound wave" was one of the earliest predictions of dusty plasma theory

- Experiments on the dust acoustic wave have been ongoing since the earliest days of dusty plasma research.
- DDWs have been studied in RF and DC glow discharge plasmas, in Q-machine plasmas, in hot filament discharge plasmas, and under microgravity conditions.
- A wide variety of experiments have been performed:
 - Experiments on self-excited and driven waves
 - Shock-driven waves

Dust acoustic waves (DAW) - I

 The original derivation of the one-dimensional dust acoustic wave was given by N. N. Rao, et al., [PSS, (1990)].

$$\omega^{2} = \frac{k^{2}C_{D}^{2}}{\left(1 + k^{2}\lambda_{D}^{2}\right)}; \quad \text{where, } C_{D} = \left[\frac{\left(\frac{T_{i}}{m_{d}}\right)\varepsilon Z^{2}}{1 + \left(\frac{T_{i}}{T_{e}}\right)\left(1 - \varepsilon Z\right)}\right]^{\frac{1}{2}} \qquad \lambda_{D}^{-2} = \lambda_{De}^{-2} + \lambda_{Di}^{-2} \quad \text{and} \quad \varepsilon = \frac{n_{d}}{n_{i}}$$

DAW/DDW basic properties - I

- Early experiments on DAWs focused on characterizing the basic properties of self-excited waves.
- The first experimental result was reported by Barkan, et al., PoP, 1995.

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He-Ne laser as the light source.

Dust acoustic waves as a diagnostic for charge

In the long λ limit, phase velocity of DAW/DDW is:

$$\frac{\omega}{k} = \left[\left(\frac{T_i}{m_d} \right) \varepsilon Z^2 \right]^{\frac{1}{2}}; \qquad \varepsilon = \frac{n_d}{n_i}$$

Use the phase velocity of the DDW, measured dust number density, and ion number density to estimate grain charge: $q_d = -Z_d e$

$$\label{eq:r_d} \begin{split} r_{\rm d} &= 0.8 \; \mu m \\ m_{\rm d} &= 6 \; x \; 10^{-16} \; \text{kg} \\ \epsilon &= 2 \; \text{to} \; 5 \; x \; 10^{-4} \\ T_{\rm i} \; (\text{est.}) &= 0.03 \; \text{eV} \\ v_{\text{phase}} &\sim 12 \; \text{cm/s} \end{split}$$

C. Thompson, et al., PoP (1997)

DUSTY PLASMA STUDIES: MAGNETIC FIELD EFFECTS

Magnetizing a dusty plasma

Magnetization criterion

- Magnetic forces will be comparable to the other forces acting upon the grain

- Challenges:
 - Dust grains charge, $Z_d \sim 1000$
 - Dust grain mass, $m_d > 10^8 m_{ion}$
- That is: $q_d/m_d \ll e/m_{ion} \ll e/m_{elec}$
- Key parameters:
 - gyroradius to exp. size:
 - gyrofreq. to collision freq.:
 - magnetic to gravitational force:

Key result: Maximize B/a

$$\frac{\rho}{L} \sim \frac{a^2 v_d}{BL} << 1$$

$$\frac{\omega_c}{v_{dn}} \sim \frac{B}{aP} > 1$$

$$\frac{F_m}{F_g} = \frac{Q_d v_d B}{m_d g} \sim \frac{v_d B}{a^2} \ge$$

Properties of a magnetized dusty plasmas

- Strong magnetic field will guide ion and electron transport to surfaces for electrodes and dust grains.
- Charging of surfaces and dust grains may be significantly modified.
- Formation of sheaths will become asymmetric parallel and perpendicular to the magnetic field.
- How is the growth of small grains affected at high magnetic field - e.g., formation of elongated particles?
- Can 2D ordering of dust be maintained at high B?

Modification of a comet's tail: an astrophysical example

- December, 2011: Comet Lovejoy passes near the Sun
- The dust tail modified by plasma flowing along magnetic field lines

Movie from NASA: http://www.nasa.gov/multimedia/videogallery/index.html?media_id=124915541

Magnetized Dusty Plasma Experiment (MDPX)

- MDPX project:
 - Develop a fully magnetized dusty plasma
 - Develop flexible, multi-configuration magnetic geometry
 - Operate as a multi-user facility
- Two primary <u>scientific questions</u>:
 - As a dusty plasma becomes magnetized how do the structural, thermal, charging, and collective properties of the system evolve?
 - If a dusty plasma has magnetic particles how does the system evolve in the presence of uniform and non-uniform magnetic fields?

Parameter space for MDPX device

	dust spacing	dust radius	charge	magnetic field	dust velocity	electron Debye length	gyroradius			
	d	а	Z	В	v	L _{de}	Dg	d/L_{de}	d/Dg	
	(m)	(µm)		(T)	(m/s)	(cm)	(m)	_		
Saturn's E-ring								0.1	2.74E-14	
Interstellar molecular clouds						7.43E+00		0.3	1.32E-12	
Star-forming region						7.43E-03		1.4	1.53E-05	
Process plasma						I.92E-04		1.5	6.91E-06	
Fusion experiment								100	I.18E-02	and the second se
MDPX - large dust	0.0002	1	4000	1	0.01	4.07E-04	I.44E-0I	0.5	1.50E-03	
	0.001	0.25	500	4	0.005	4.075.04	2 25E 02	2.5	4 45E 01	
ribra - smail dust	0.001	0.23	500	4	0.005	4.V/ E=V4	2.23E-03	2.5	4.43E-VI	

n ~ 10¹⁵ m⁻³ T_e ~ 3 eV

Not to scale

AUPSL Adapted from Mendis and Rosenberg, Ann. Rev. Astron. Astrophys., 32, 419 (1994)

MDPX superconducting magnet system

- Four superconducting coils in a split-bore cryostat
- 1.3 mm diameter, copper wire with Nb-Ti filaments
- Use of cryogen-free cooling system (i.e., no liquid helium)

MDPX parameters: final "as-built" system

Magnetic field: Magnetic field gradient: Magnet cryostat: Magnet material: 3.0 T (to date); 4 T (max) I - 2 T /m 50 cm ID / 127 cm OD NbTi superconductor; cryogen-free system

Magnet performance:

Experiment volume:* Uniform region:* Total cold mass: T_{critical} = 6.4 K; T_{minimum} = 4.2 - 4.8 K 45 cm dia. x 175 cm axial 20 cm dia. x 20. cm axial ~2.5 tons

Self-organization in dusty plasmas

 $\exp(-r / \lambda_{D})$

- Γ (coupling parameter) is indicative of the selforganizing, emergent properties of dusty plasmas.
- A dusty plasma can be used as a model system to investigate problems in soft-matter physics.
- Assume dust particles interact via a screened Coulomb interaction (Yukawa, Debye-Hückel):

Imposed order in a dusty plasma at high magnetic field: experiment setup

Lower electrode: powered Upper electrode: electrically floating, wire mesh Electrode gap: 62 mm

Dust cloud: 30 - 40 mm below upper electrode

Key scale lengths: (B = I T, P = I20 mTorr) $\lambda_{mfp-ion} \sim 0.6 \text{ mm}$ $\lambda_{mfp-electron} \sim 8 \text{ mm}$ $\lambda_{De} \sim 0.2 \text{ mm}$ $r_{Le} \sim 0.006 \text{ mm}$ $r_{Li} \sim 0.15 \text{ mm}$

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Imposed order in a dusty plasma at high magnetic field: video of dust motion

Dust motion: p = 145 mTorr 2 micron particles B = 2 T

Sum over 100 frames and image contrast enhanced.

Note the grid-like pattern followed by the dust grains (in white).

There are also filaments (light gray areas) beginning to form.

Dust grid structure becomes degraded with higher neutral pressure

- Comparison of grid formation: 0.5 micron particles, B = 1.5 T Sum over 100 frames and image contrast enhanced
- The ordering of the dust into the grid pattern decreases with higher pressure.
- Fluid-like motion increases with higher pressure.
- Note differences in plasma glow and filamentation with pressure. (lower p, more filaments)

Structure persists over a range of magnetic field, pressure, and particles sizes

Close-up of a 450 x 450 pixel region

d = 0.61 mm

Close-up of a 125 x 125 pixel region

2 T 145 mTorr 2 μm d = 0.65 mm

I.5 T 53 mTorr 0.5 μm

d = 0.69 mm

Grid pattern in dust: ~0.65 ± 0.02 mm

Mesh wire used: **#40 titanium mesh** Wire: 0.25 mm diameter Center-to-center spacing: 0.635 mm

Beginning studies of dust acoustic/dust density waves at high magnetic field

Most recent wave measurements: B = 1.0 T2 micron diameter particles p = 140 mTorr

Original video: 100 fps, slowed by 1/10th

Beginning studies of dust acoustic/dust density waves

Most recent wave measurements: B = 1.0 T 2 micron diameter particles p = 140 mTorr

- Analysis of global particle motion using particle image velocimetry (PIV).
- Plots show a distribution of the measured PIV velocity vectors as a function of zposition (i.e., into the plane).
- This is suggestive of the circulation of the particles from the front of the cloud $(+V_x)$ to the back of the cloud $(-V_x)$.

Summary: current status of the MDPX device

- Engineering status:
 - MDPX is generally performing well; peak field B = 3 T
 - Operations are shifting from a testing to operational mode.
 - Gaining experience with making reliable and reproducible plasmas and dusty plasmas.
- <u>Research status:</u>
 - Experiments are focused on the interaction of magnetized ions and electrons with charged dust.
 - Use of wire mesh has led to the observation of a new of "imposed ordering" in a dusty plasma whose characteristics strongly differ from previously observed plasma crystals.
 - Thomas, et al., J. Plasma Phys., 81, 345810206 (2015); Thomas, et al., *Phys. Plasmas*, 22, 030701 (2015).
 - We are making detailed observations of transport and instabilities.
 - Seeking to expand collaborations with national and international groups.

OUTLOOK FOR DUSTY PLASMA RESEACH

Outlook for dusty plasma research is very promising

- Dusty plasma as "model systems" for soft condensed matter, fluid systems and statistical mechanics - need new insights and people to help make these connections.
- A unified model of dust grain charging in plasmas still remains elusive can a model be developed that works for lab, fusion, and space plasmas?
- Upcoming space missions to Jupiter, comets, Moon, Mars will involve study of charged dust or charged ice in solar system environment - need a new generation of lab studies to support these missions.
- Several groups around the world are studying magnetic field effects need new models, theories, and diagnostic tools to understand experimental observations.
- Dust formation and control in fusion is a major issue for plasma-material interaction studies dust is a contaminant, but can it be used to control and fuel the plasma or forn disruption mitigation?
- New "multi-user" dusty plasma lab facilities for ground- and space-based research are coming online.

References

- Journals
 - TPS IEEE Transactions on Plasma Science
 - PoP Physics of Plasmas
 - PRL Physical Review Letters
 - PRE Physical Review E
 - PPCF Plasma Physics and Controlled Fusion
 - PSS Planetary and Space Science
- Textbooks
 - Introduction to Dusty Plasma Physics P. Shukla and A. Mamun
 - Physics and Applications of Complex Plasmas S.Vladimirov, K. Ostrikov, and A. Samarian
 - Plasma Physics A. Piel

Selected list of institutions involved in dusty plasma research

• US

- Auburn University (Physics)
- Baylor University (Physics)
- Caltech (Physics)
- University of Alabama at Huntsville (Mech. Eng.)
- University of California San Diego (Elec. Eng.)
- University of Colorado (Physics)
- University of Iowa (Physics)
- University of Michigan (Elec. Eng.)
- University of Minnesota (Mech. Eng.)
- MIT (Nucl. Eng.)
- Virginia Tech
- Wittenberg University
- Los Alamos National Lab
- Princeton Plasma Physics Lab
- Naval Research Lab
- International
 - Germany: U. Kiel, U. Giessen, U. Greifswald, U. Bochum, Germany Aerospace Center (DLR)
 - France: CNRS Marseilles, U. Orleans
 - Sweden: Royal Institute of Technology, Univ. of Stockholm
 - Japan: U. Kyoto,
 - India: Inst. Plasma Research (IPR), U. Delhi



